

Lower bounds for blow up time of the p -Laplacian equation with damping term

YAVUZ DİNÇ, ERHAN PIŞKİN, CEMİL TUNÇ

ABSTRACT. In this work deals with the p -Laplacian wave equation with damping terms in a bounded domain. Under suitable conditions, we obtain a lower bounds for the blow up time. Our result extends the recent results obtained by Baghaei (2017) and Zhou (2015), for $p > 2$.

1. INTRODUCTION

In this work, we study the following p -Laplacian equation with strong and weak damping terms

$$(1) \quad \begin{cases} u_{tt} - \operatorname{div} \left(|\nabla u|^{p-2} \nabla u \right) - a \Delta u_t + b u_t = |u|^{q-2} u, & x \in \Omega, \quad t > 0, \\ u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), & x \in \Omega, \\ u(x, t) = 0, & x \in \partial\Omega, \quad t > 0, \end{cases}$$

where $\Omega \subset \mathbb{R}^n$ ($n = 2, 3, \dots$) is a bounded domain with a smooth boundary $\partial\Omega$; and $u_0(x) \in W_0^{1,p}(\Omega)$, $u_1(x) \in L^2(\Omega)$, $p > 2$. $a \geq 0$, $b > -a\rho_1$ with $\rho_1 > 0$ is the first eigenvalue of the operator $-\Delta$ under homogeneous Dirichlet boundary conditions and

$$(2) \quad \begin{cases} 2 < q < \infty & \text{if } n = 2, \\ 2 < q \leq \begin{cases} \frac{2n}{n-2}, & \text{for } a > 0 \\ \frac{2n-2}{n-2}, & \text{for } a = 0 \end{cases} & \text{if } n \geq 3. \end{cases}$$

When $p = 2$, (1) is reduced to the following wave equation

$$(3) \quad u_{tt} - \Delta u - a \Delta u_t + b u_t = |u|^{q-2} u.$$

In 2006, Gazzalo and Squassina [2] studied problem (3). They proved the local existence, global existence and blow up of solutions. Later, some authors studied the lower bounds for the blow up time under some conditions,

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see [1, 7, 9]. Also, in recent years, some authors investigate the lower bounds for blow up time for hyperbolic type equations, see [3, 4, 6, 8].

Inspired by the above papers, in this paper we consider the lower bound for the blow up time of solutions (1). Our result improves the recent results obtained by Baghaei [1] and Zhou [9], for $p > 2$.

Now, for the problem (1), we define the functionals, the potential well depth d and the unstable set U , are given as

$$\begin{aligned}
 I(u) &= \|\nabla u\|_p^p - \|u\|_q^q, \\
 J(u) &= \frac{1}{p} \|\nabla u\|_p^p - \frac{1}{q} \|u\|_q^q, \\
 E(t) &:= E(u, u_t) = J(u) + \frac{1}{2} \|u_t\|^2, \\
 (4) \quad d &= \inf_{u \in H_0^1(\Omega) \setminus \{0\}} \max_{\lambda \geq 0} J(\lambda u), \\
 U &= \left\{ u \in W_0^{1,p} : J(u) \leq d \text{ and } I(u) < 0 \right\}
 \end{aligned}$$

with above assumptions, the results on the existence of local solutions and the nonexistence of the solutions of the problem (1) can be reformulated as follows:

- (i) Assume that q satisfy (2), then there exist $T > 0$ and a unique solution u problem (1) satisfying

$$\begin{aligned}
 u &\in C\left([0, T], W_0^{1,p}(\Omega)\right) \cap C^1\left([0, T], L^2(\Omega)\right) \cap C^2\left([0, T], H^{-1}(\Omega)\right), \\
 u_t &\in L^2\left((0, T), H_0^1(\Omega)\right).
 \end{aligned}$$

- (ii) The solution u of problem (1) blows up at a finite time T if and only if there exists $\bar{t} \in [0, T)$ such that $u(\bar{t}) \in U$ and $E(u(\bar{t}), u_t(\bar{t})) \leq d$.

$$\begin{aligned}
 (5) \quad E(t) &= \frac{1}{2} \|u_t\|^2 + \frac{1}{p} \|\nabla u\|_p^p - \frac{1}{q} \|u\|_q^q, \\
 E(0) &= \frac{1}{2} \|u_1\|^2 + \frac{1}{p} \|\nabla u_0\|_p^p - \frac{1}{q} \|u_0\|_q^q.
 \end{aligned}$$

Lemma 1 ([5]). *Let $\Omega \subset R^n$ ($n \geq 2$). Let u be a non-negative piecewise C^1 function defined in Ω with $u(x) = 0, x \in \partial\Omega$. So*

$$(6) \quad \int_{\Omega} |u|^{2s} \, dx \leq \delta \left(\int_{\Omega} |\nabla u|^2 \, dx \right)^s$$

satisfies for

$$\begin{cases} s > 1, & \text{if } n = 2, \\ 1 < s < \frac{n}{n-2}, & \text{if } n \geq 3, \end{cases}$$

with $\delta = \left(\frac{n-1}{n^{\frac{3}{2}}}\right)^{2s} |\Omega|^{1-\frac{(n-2)}{n}s}$.

Theorem 1. Assume that $q > \frac{p+2}{2}$ holds. Let $u(x, t)$ be the solution of the problem (1), which blows up at a finite time T^* . Then

$$\int_{\psi(0)}^{\infty} \frac{d\tau}{qM + \tau + \varepsilon q p^{\frac{2(q-1)}{p}} 2^{\frac{2(q-1)}{p}-2} M^{\frac{2(q-1)}{p}} + \varepsilon p^{\frac{2(q-1)}{p}} q^{\frac{p-2q+2}{p}} 2^{\frac{2(q-1)}{p}-2} \tau^{\frac{2(q-1)}{p}}} \leq T^*,$$

where

$$(7) \quad \begin{cases} \psi(0) = \|u_0\|_q^q, \\ \varepsilon = \left(\frac{n-1}{\frac{n}{2}}\right)^{2(q-1)} |\Omega|^{1-\frac{(n-2)(q-1)}{n}}, \\ M = \frac{1}{2} \|u_1\|^2 + \frac{1}{p} \|\nabla u_0\|_p^p - \frac{1}{q} \|u_0\|_q^q. \end{cases}$$

Proof. Multiplying the equation of (1) by u_t , then integrating the result over Ω , we have

$$\begin{aligned} \int_{\Omega} u_t \left[u_{tt} - \operatorname{div} \left(|\nabla u|^{p-2} \nabla u \right) - a \Delta u_t + b u_t \right] dx &= \int_{\Omega} u_t |u|^{q-2} u dx, \\ \frac{1}{2} \frac{d}{dt} \int_{\Omega} |u_t|^2 dx + \frac{1}{p} \frac{d}{dt} \int_{\Omega} |\nabla u|^p dx + a \int_{\Omega} |\nabla u_t|^2 dx + b \int_{\Omega} |u_t|^2 dx \\ &= \frac{1}{q} \frac{d}{dt} \int_{\Omega} |u|^q dx. \end{aligned}$$

Since

$$\frac{d}{dt} \left(\frac{1}{2} \|u_t\|^2 + \frac{1}{p} \|\nabla u\|_p^p - \frac{1}{q} \|u\|_q^q \right) = -a \|\nabla u_t\|^2 - b \|u_t\|^2,$$

the above calculations imply $E'(t) \leq 0$, with

$$E(t) := \frac{1}{2} \|u_t\|^2 + \frac{1}{p} \|\nabla u\|_p^p - \frac{1}{q} \|u\|_q^q.$$

So,

$$(8) \quad E(t) \leq E(0) = M,$$

where M is defined in (7) and

$$(9) \quad \|u_t\|^2 + \frac{2}{p} \|\nabla u\|_p^p = 2E(t) + \frac{2}{q} \|u\|_q^q \leq 2M + \frac{2}{q} \|u\|_q^q.$$

Now define

$$\psi(t) = \|u\|_q^q,$$

then thanks to Cauchy's and embedding ($L^p(\Omega) \hookrightarrow L^2(\Omega)$, $p > 2$) inequalities and Lemma 1 with $s = q - 1$ and $\delta = \varepsilon$, we obtain

$$\begin{aligned} \psi'(t) &= q \int_{\Omega} |u|^{q-2} u u_t dx \\ &\leq \frac{q}{2} \left(\|u_t\|^2 + \|u\|_{\frac{2(q-1)}{2(q-1)}}^{2(q-1)} \right) \\ &\leq \frac{q}{2} \left(\|u_t\|^2 + \varepsilon \|\nabla u\|_2^{2(q-1)} \right) \end{aligned}$$

$$\begin{aligned}
&\leq \frac{q}{2} \left(\|u_t\|^2 + \varepsilon \|\nabla u\|_p^{2(q-1)} \right) \\
&= \frac{q}{2} \left(\|u_t\|^2 + \varepsilon \left(\|\nabla u\|_p^p \right)^{\frac{2(q-1)}{p}} \right).
\end{aligned}$$

By the (9) and definition of ψ , we get

$$\begin{aligned}
\psi'(t) &\leq \frac{q}{2} \left[2M + \frac{2}{q}\psi(t) + \varepsilon \left(\frac{p}{2}\right)^{\frac{2(q-1)}{p}} \left(2M + \frac{2}{q}\psi(t) \right)^{\frac{2(q-1)}{p}} \right] \\
&\leq \frac{q}{2} \left[\begin{array}{c} 2M + \frac{2}{q}\psi(t) \\ + \varepsilon \left(\frac{p}{2}\right)^{\frac{2(q-1)}{p}} 2^{\frac{2(q-1)}{p}-1} \left((2M)^{\frac{2(q-1)}{p}} + \left(\frac{2}{q}\psi(t)\right)^{\frac{2(q-1)}{p}} \right) \end{array} \right] \\
&= \frac{q}{2} \left[2M + \frac{2}{q}\psi(t) + \varepsilon \frac{p^{\frac{2(q-1)}{p}}}{2} \left((2M)^{\frac{2(q-1)}{p}} + \left(\frac{2}{q}\psi(t)\right)^{\frac{2(q-1)}{p}} \right) \right] \\
&= qM + \psi(t) + \varepsilon q \frac{p^{\frac{2(q-1)}{p}}}{4} \left(2^{\frac{2(q-1)}{p}} M^{\frac{2(q-1)}{p}} + \frac{2^{\frac{2(q-1)}{p}}}{q} \psi^{\frac{2(q-1)}{p}}(t) \right) \\
&= qM + \psi(t) + \varepsilon qp^{\frac{2(q-1)}{p}} 2^{\frac{2(q-1)}{p}-2} M^{\frac{2(q-1)}{p}} \\
&\quad + \varepsilon p^{\frac{2(q-1)}{p}} q^{\frac{p-2q+2}{p}} 2^{\frac{2(q-1)}{p}-2} \psi^{\frac{2(q-1)}{p}}(t),
\end{aligned}$$

where used $(a+b)^\varpi \leq 2^{\varpi-1}(a^\varpi + b^\varpi)$. Thus, the above inequality implies that

$$\begin{aligned}
\frac{d\psi(t)}{dt} &\leq qM + \psi(t) + \varepsilon qp^{\frac{2(q-1)}{p}} 2^{\frac{2(q-1)}{p}-2} M^{\frac{2(q-1)}{p}} \\
(10) \quad &\quad + \varepsilon p^{\frac{2(q-1)}{p}} q^{\frac{p-2q+2}{p}} 2^{\frac{2(q-1)}{p}-2} \psi^{\frac{2(q-1)}{p}}(t).
\end{aligned}$$

Since $\lim_{t \rightarrow T^*} \psi(t) = \infty$, we get from (10)

$$\int_{\psi(0)}^{\infty} \frac{d\tau}{qM + \tau + \varepsilon qp^{\frac{2(q-1)}{p}} 2^{\frac{2(q-1)}{p}-2} M^{\frac{2(q-1)}{p}} + \varepsilon p^{\frac{2(q-1)}{p}} q^{\frac{p-2q+2}{p}} 2^{\frac{2(q-1)}{p}-2} \tau^{\frac{2(q-1)}{p}}} \leq T^*.$$

This completes the proof of the main theorem. \square

2. CONCLUSION

In this work, we obtained the lower bounds for blow up time of the non-linear p -Laplacian equation with damping terms in a bounded domain. This improves and extends many results in the literature.

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YAVUZ DİNÇ

DEPARTMENT OF ELECTRONICS
TECHNICAL MARDİN ARTUKLU UNIVERSITY
MARDİN
TURKEY
E-mail address: yavuzdinc@artuklu.edu.tr

ERHAN PIŞKİN

DEPARTMENT OF MATHEMATICS
DİCLE UNIVERSITY
21280, DIYARBAKIR
TURKEY
E-mail address: episkin@dicle.edu.tr

CEMİL TUNÇ

DEPARTMENT OF MATHEMATICS
VAN YÜZÜNCÜ YIL UNIVERSITY
VAN
TURKEY
E-mail address: cemtunc@yahoo.com